

Tilted-axis-cranking covariant density functional theory for high-spin spectroscopy of ^{69}Ga *

Y. P. Wang (王一平)[✉] Y. K. Wang (王亚坤)[✉] P. W. Zhao (赵鹏巍)[†][✉]

State Key Laboratory of Nuclear Physics and Technology, School of Physics, Peking University, Beijing 100871, China

Abstract: The tilted-axis-cranking covariant density functional theory is applied to investigate the three newly-observed positive-parity bands SI, SII, and SIII in ^{69}Ga . The energy spectra and angular momenta are calculated, and they agree closely with experimental data. For band SI, pairing correlations are crucial for the states with spin $I \leq 23/2\hbar$. Bands SII and SIII are suggested to be signature partner bands with positive and negative signatures, respectively. By analyzing the angular momentum alignments, we reveal that the $g_{9/2}$ protons and neutrons are crucial in the collective structures of ^{69}Ga . The transition probabilities $B(E2)$ for these bands are predicted, awaiting further experimental verification.

Keywords: high-spin spectroscopy, collective structure, cranking covariant density functional theory

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I. INTRODUCTION

The spectroscopic properties of nuclei in the $A \approx 70$ mass region have attracted considerable interest for experimental and theoretical studies over the past few decades. They provide an ideal laboratory for studying the interplay between the single-nucleon and collective motions. Generally, at low spins, the properties of the $A \approx 70$ nuclei are governed by single-particle excitations involving low- j orbitals, such as $1f_{5/2}$, $2p_{3/2}$, and $2p_{1/2}$. At intermediate and high spins, the high- j deformation-driving $1g_{9/2}$ orbitals come into effect and induce the collective behavior. The interplay between the single-nucleon and collective motions in this mass region generates many interesting phenomena, such as superdeformed bands [1–3] and magnetic rotations [4, 5].

Compared with the widely studied even-even nuclei [6–12], studies on the spectroscopic properties of the odd-mass nuclei with $A \approx 70$ are rare. The spectroscopic properties of odd-mass nuclei are worth exploring because the existence of the unpaired nucleon would result in additional complexity and more interesting physics. Available studies on the odd-mass nuclei in the $A \approx 70$ mass region have already provided valuable insights into the evolution of the single-particle states around the $Z = 28$ closed shell [13, 14] and the effects of collective motions [15, 16].

With three protons beyond the $Z = 28$ closed shell, the

odd-mass Ga isotopes present diverse spectral structures and has captured much interest [17–20]. Many experimental efforts have been devoted to measuring the spectra of ^{69}Ga since 1956 [21]. These investigations reported various low-lying negative-parity states generated by single-particle excitations involving $1f_{5/2}$, $2p_{3/2}$, $2p_{1/2}$, and $1f_{7/2}$ orbitals [22–26]. Very recently, through the $^{26}\text{Mg}(^{48}\text{Ca}, p4n\gamma)$ multi-nucleon transfer reaction, the negative-parity level scheme of ^{69}Ga has been considerably extended. In addition, three positive-parity bands, namely bands SI, SII, and SIII, have been observed [27], and further theoretical studies are required for better understanding their nature.

Theoretically, the spectroscopic properties of nuclei in the $A \approx 70$ mass region have been investigated using various approaches, *e.g.*, the interacting boson-fermion model [28, 29], configuration-interaction shell model [30, 31], and tilted-axis-cranking covariant density functional theory (TAC-CDFT) [32–34]. The TAC-CDFT [35, 36] starts from a universal density functional; thus, it can provide a microscopic understanding on the nuclear single-nucleon and collective motions without introducing adjustable parameters. Moreover, the TAC-CDFT exploits the Lorentz symmetry, and thus, nuclear currents that are important for the description of nuclear rotations [35, 37–39] can be considered self-consistently.

In this study, the TAC-CDFT is applied to investigate the three newly-observed positive-parity bands SI,

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[†] E-mail: pwzhao@pku.edu.cn

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SII, and SIII in ^{69}Ga . A microscopic interpretation for the collective structures in ^{69}Ga is provided.

II. THEORETICAL FRAMEWORK

The detailed formalism of the CDFT and TAC-CDFT is given in Refs. [32, 35, 36, 40]. Here, a brief introduction is presented.

The starting point of the CDFT is a universal density functional [41–44]. To perform the TAC calculations based on the CDFT, we transform the functional into a body-fixed frame rotating with a constant angular velocity ω , whose direction is determined by minimizing self-consistently the total energy in the rotating frame. The relativistic Kohn-Sham equation in the rotating body-fixed frame is expressed as

$$[\alpha \cdot (-i\nabla - \mathbf{V}) + \beta(m + S) + V - \omega \cdot \hat{\mathbf{J}}] \psi_k = \varepsilon_k \psi_k, \quad (1)$$

where $\hat{\mathbf{J}}$ is the nuclear total angular momentum. The scalar field S and vector field V^μ are connected in a self-consistent manner to the nucleon densities and currents [32]. The iterative solution of Eq. (1) yields single-particle states and single-particle energies. The obtained single-particle states are then used to calculate the expectation values of three components $\langle J_i \rangle$ of the angular momentum, total binding energies, $B(E2)$ transition probabilities, etc. The magnitude of ω is connected to the total angular momentum quantum number I by the semiclassical relation $\langle \hat{\mathbf{J}} \rangle \cdot \langle \hat{\mathbf{J}} \rangle = I(I + 1)$.

In the TAC-CDFT framework, the pairing correlations are often treated using the Bogoliubov method [42, 44] and the shell-model-like approach (SLAP) [45]. In comparison, the SLAP treats the pairing correlations exactly by diagonalizing the many-body Hamiltonian in a properly truncated many-particle configuration (MPC) space. Thus, it can conserve the particle number and avoid the pairing collapse; for details, see Refs. [45–48].

Here, to investigate the collective structures in ^{69}Ga , we solve the relativistic Kohn-Sham equation using a three-dimensional harmonic-oscillator basis in Cartesian coordinates with ten major shells, which ensures convergence for nuclei in the $A \approx 70$ mass region [32]. The PC-PK1 relativistic density functional [49] is adopted. Such a functional has demonstrated a high predictive power in describing nuclear masses [50, 51] and shapes [52–54], magnetic and antimagnetic rotations [32, 36, 55–57], chiral rotation [58–61], and nuclear fission [62, 63].

III. RESULTS AND DISCUSSIONS

To identify suitable configurations for the positive-parity bands SI, SII, and SIII, we perform cranking calculations using the configuration constrained approach [35]

without considering the pairing correlations. By comparing the obtained results with the experimental data, we identify the two most likely configurations. Their energies and quadrupole deformation parameters β and γ at rotational frequency $\hbar\omega = 0$ MeV are listed in Table 1. The configuration $\nu(g_{9/2})^4(pf)^6 \otimes \pi(g_{9/2})^1(pf)^2$ is assigned to band SI. The configuration $\nu(g_{9/2})^4(pf)^6 \otimes \pi(g_{9/2})^3$ with positive and negative signatures is assigned to bands SII and SIII, respectively. Both configurations involve several valence $g_{9/2}$ protons and neutrons that would induce nuclear collective rotations. The γ deformations for these two configurations are zero, and the corresponding β deformations are large. The particle nature of these configurations generates principal axis rotations, which are consistent with the $\Delta I = 2$ character of the bands under consideration.

The calculated energy spectra and angular momenta for the positive-parity band SI are shown in Fig. 1, in comparison with experimental data [27]. Without considering the pairing correlations, the theoretical results provide a good description for the states with spin $I > 23/2\hbar$. Nevertheless, the observed excitation energies and total angular momenta for the states with $I \leq 23/2\hbar$ are overestimated, which might be related to the lack of pairing correlations. To better clarify this point, we perform calculations with pairing correlations considered by the SLAP. In the calculations, the monopole pairing force is adopted. The effective neutron and proton pairing strengths are 0.7 MeV and 0.75 MeV, respectively, determined by reproducing the experimental odd-even mass differences. The MPC space is truncated at an excitation energy cutoff [64] of $E_c = 10$ MeV. A larger MPC space with a renormalized pairing strength essentially yields the same results, indicating the convergence of the MPC space. We observe that the calculated energy spectra and angular momenta are in better agreement with the experimental data after the inclusion of the pairing correlations. Note that, for the $I \geq 29/2\hbar$ region, two $g_{9/2}$ neutrons can be easily excited to the negative-parity $2p_{3/2}$ orbitals, which leads to an oblate configuration labeled as $\nu(g_{9/2})^2(pf)^8 \otimes \pi(g_{9/2})^1(pf)^2$. Therefore, configuration constrained calculations must be conducted to maintain the same configuration as that for the low-spin region.

The calculated energy spectra and angular momenta for the positive-parity bands SII and SIII in ^{69}Ga , as well as their comparison with the data [27], are shown in Fig. 2. The spins for the band heads of SII and SIII are

Table 1. Calculated total energies, deformations β and γ , and assigned configurations for positive-parity bands SI–SIII in ^{69}Ga .

Band	E / MeV	β	γ	Configuration
SI	−596.2	0.41	0.0°	$\nu(g_{9/2})^4(pf)^6 \otimes \pi(g_{9/2})^1(pf)^2$
SII, SIII	−595.6	0.57	0.0°	$\nu(g_{9/2})^4(pf)^6 \otimes \pi(g_{9/2})^3$

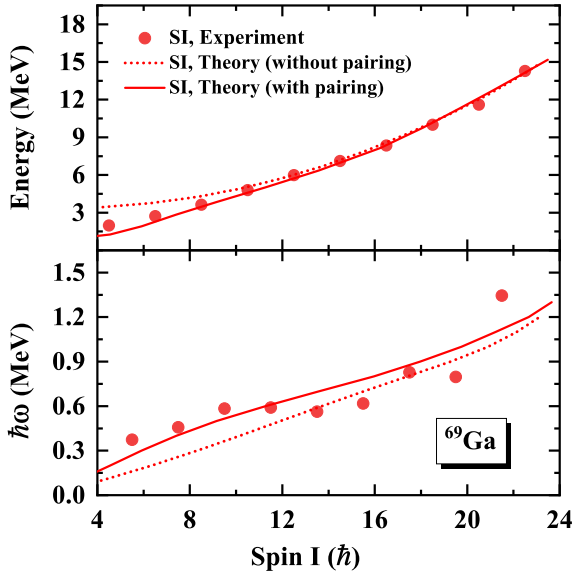


Fig. 1. (color online) Calculated excitation energies (upper panel) and rotational frequency (lower panel) as functions of the angular momenta for the positive-parity band SI in ^{69}Ga with and without pairing, in comparison with the data [27].

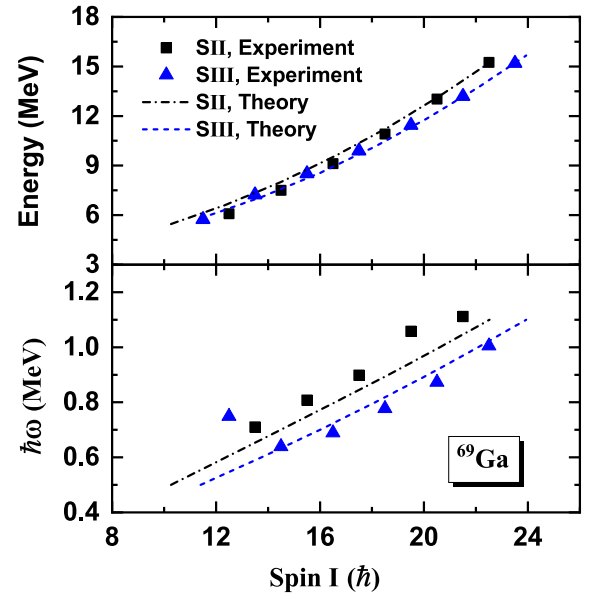


Fig. 2. (color online) Calculated excitation energies (upper panel) and rotational frequency (lower panel) as functions of the angular momenta for the positive-parity bands SII and SIII in ^{69}Ga , in comparison with the data.

already larger than $23/2\hbar$, making the pairing correlations negligible for these two bands. The experimental energies and angular momenta for both SII and SIII are properly reproduced by the present calculations. Since these two bands are based on the same configuration with opposite signatures, they are suggested as signature partner bands, in which the one with negative signature is more energy favored. A signature inversion in bands SII and SIII occurs at $35/2\hbar$, which cannot be reproduced by the present calculation, in which these two bands are calculated independently based on two different configurations. The interactions between these two configurations may be important for reproducing such a signature inversion phenomenon. Beyond-mean-field methods, such as relativistic configuration-interaction density functional theory (ReCD) [65–68], may be useful in resolving the discrepancy between the calculated results and data.

As mentioned earlier, all the studied bands involve nucleons in $g_{9/2}$ orbitals, and the reasonable description of these bands gives us confidence to provide a microscopic interpretation of the role of the $g_{9/2}$ orbitals in the collective structures in ^{69}Ga . In the present microscopic calculations, the nuclear total angular momenta are generated from all the individual nucleons inside a self-consistent meanfield potential. Band SI has four $g_{9/2}$ valence neutrons and one $g_{9/2}$ valence proton. Bands SII and SIII have four $g_{9/2}$ valence neutrons and three $g_{9/2}$ valence protons. The angular momentum alignments of these $g_{9/2}$ nucleons calculated using the TAC-CDFT without the pairing correlations are shown in Fig. 3. For band SI, as shown in Fig. 3 (a), the contribution from the four neutrons in the $g_{9/2}$ orbitals increases with the rotational fre-

quency ω and is significant in generating the total angular momentum. In contrast, the alignment of the unpaired $g_{9/2}$ proton is almost constant. For band SII, as shown in Fig. 3 (b), the alignments of the valence $g_{9/2}$ neutrons and protons increase simultaneously with ω , and their contributions to the total angular momentum are about the same order. The behavior of band SIII is similar to that of its signature partner band SII and is shown in Fig. 3 (c).

Finally, the $B(E2)$ values for the positive-parity bands SI, SII, and SIII are predicted and are shown as functions of the angular momenta in Fig. 4. The $B(E2)$ values exhibit a decreasing tendency for all these bands. This is caused by the fact that their deformation parameters β decrease with the total angular momenta. For band SI, when the pairing correlations are considered, the $B(E2)$ values

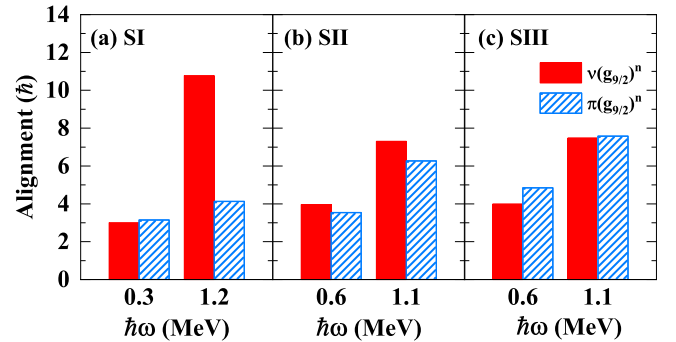


Fig. 3. (color online) Angular momentum alignments of the $g_{9/2}$ proton and neutron particles for the bands SI (a), SII (b) and SIII (c), calculated using the TAC-CDFT without the pairing correlations.

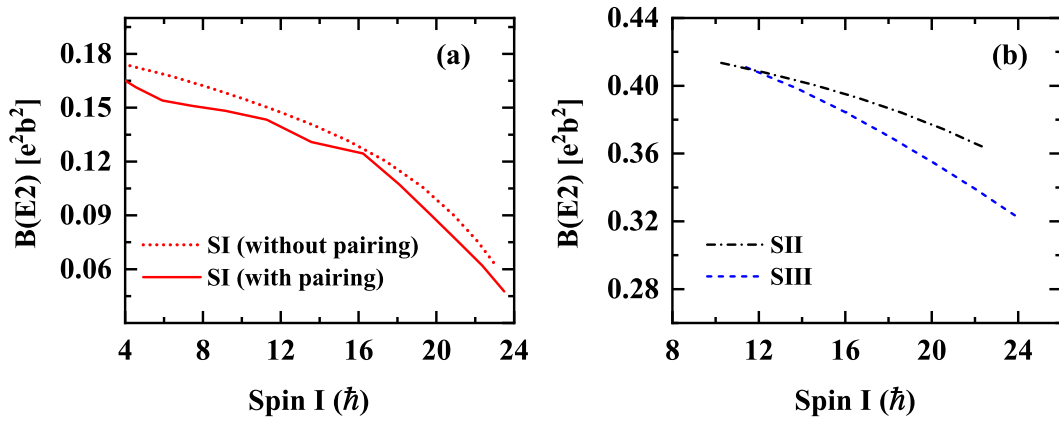


Fig. 4. (color online) Calculated $B(E2)$ values as functions of the angular momenta for the positive-parity bands (a) SI and (b) SII and SIII.

decrease slightly. For the signature partner bands SII and SIII, the calculated $B(E2)$ values are close at spin $I \approx 23/2\hbar$. However, at higher spins, the $B(E2)$ values of band SII become larger than those of band SIII. Future experimental efforts that incorporate the measurement of lifetimes of these states are required to validate our predictions.

IV. SUMMARY

The tilted-axis- cranking covariant density functional theory is applied to investigate the three newly-observed positive-parity bands SI, SII, and SIII in ^{69}Ga . The energy spectra and angular momenta are calculated and compared with experimental data. For band SI, pairing correlations play a crucial role for the states with spin $I \leq 23/2\hbar$. Bands SII and SIII are suggested to be signa-

ture partner bands with positive and negative signatures, respectively. By analyzing the angular momentum alignments, we observe that the $g_{9/2}$ protons and neutrons play an important role in the collective structures of ^{69}Ga . For band SI, the two neutrons in the $g_{9/2}$ orbitals significantly contribute to the generation of the total angular momentum. For bands SII and SIII, the alignments of the valence $g_{9/2}$ neutrons and protons increase simultaneously with ω , and their contributions to the total angular momentum are about the same order. The transition probabilities $B(E2)$ for these bands are predicted and await further experimental verification.

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