

Disentangling the nature of resonances in coupled-channel models ^{*}

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Abstract: We present several possible hadronic states found in coupled-channel models within the on-shell approximation. The interaction potential is constructed as a sum of the tree-level Feynman diagrams calculated with the effective Lagrangians. Based on the recent empirical data, we illustrate the possible existence of several baryonic and mesonic states with definite quantum numbers in the model. We give their properties for the purpose of further study and discuss the potential of finding them in future experiments.

Key words: partial wave analysis, coupled-channel model, exotic state

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1 Introduction

Nowadays, Quantum Chromodynamics (QCD) has not yet been numerically resolved at the range of low energies, if not mentioning lattice QCD for the moment. Instead, in order to reveal the nature of hadronic states, a lot of sophisticated phenomenological models are developed to separate the resonance from the overlapped background and coupled-channel effect. In this energy range of the non-perturbative Quantum Chromodynamics, the coupled-channel effect makes the analysis complicated and partial wave analysis (PWA) seems to be a useful tool to analyze the baryonic and mesonic states. On the experimental side, the scattering of stable particles and the hadronic decay of heavy quarkonium are usually used to extract empirical information of resonances. The typical examples are the πN , γN [1–5], eN and pN reactions [6–13] and the decay of the charmonium states, e.g. J/ψ , ψ' and χ_{cJ} [14–16].

The analyses of the hadronic decay of heavy quarkonium are simplified by the isospin filter which only allows the contribution of isospin zero. A lot of progress has been made in this direction in recent years [17, 18]. The states with isospin one should be studied in other reactions, but they are usually convoluted with isospin zero components. In the pN reactions [6–13], the coupled-channel effects are complicated due to three or more particles in final states so the study is still in the initial stage. For the photo-induced reactions, though the non-

resonant background is large due to the enhancement of Born terms resulting from the gauge invariance [2], the unprecedented development of coupled-channel models has reliably dissociated the resonances from the background and investigated the properties of many states by treating the amount of πN and γN data on an equal footing. One kind of coupled-channel model proceeded from the Bethe-Salpeter equation in the on-shell approximation [1, 2]:

$$T = V + VGT = \frac{V}{1 - VG}, \quad (1)$$

where T stands for the total transition amplitudes and V is the interaction kernel in the tree level, respectively. The propagators matrix G is diagonal. On the basis of fundamental scattering theory, the wave function could be expanded in the eigenstates of the angular momentum. As a result, the transition amplitudes in the above equation could be decomposed into a series of partial waves with definite quantum numbers. Then the scattering equation can be solved for each partial wave separately. This PWA has proved to be a powerful tool to compare different models and experimental approaches with each other. The partial waves in different models give a direct comparison in the details of various approaches and serve as a guideline for future improvements. In addition, PWA can be used to decompose the measured sample of events into the constituent contributions of partial waves, and so give a straight link to

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the model calculation and experiment. In the following text we would like to show that several baryonic and mesonic states could be found in this scheme based on the experimental data.

2 New baryonic states in photo-induced reactions

The Bethe-Salpeter equation Eq. (1) is extraordinarily complicated because its denominator involves four-dimensional loop integration. In order to reduce this daunting integration task and numerically solve the Eq. (1) in meson-baryon interaction channels, the approximation is unavoidable. The propagator G in Eq. (1) could be split into two constituents containing the real and imaginary parts. In the so-called K -matrix approximations, the real part of the propagator is neglected. This is justified considering that the real part of the intermediate loop integrals only results in a non-observable renormalization of the coupling constants and masses of the involved particles. The validity of this approximation has been extensively tested and explored in the literature [1, 2]. After some algebraic manipulation, the scattering T -matrix is reduced to the set of equations for each partial wave [1, 2]:

$$T_{\text{fi}}^{J\pm,I} = \left[\frac{K^{J\pm,I}}{1 - iK^{J\pm,I}} \right]_{\text{fi}}, \quad (2)$$

where $J\pm$ and I are total spin, parity and isospin of the initial and final states. Below the center of mass energy 2000 MeV, the $f, i = \gamma N, \pi N, 2\pi N, \eta N, \omega N, K\Lambda$ and $K\Sigma$ channels are opened and should be included in the model. The kernel $K=V$ could be dynamically built as the $s, u,$ and t - channels with the effective interaction Lagrangians respecting chiral symmetry in the low-energy regime [1]. Recently we analyzed the current data of $K\Sigma$ photoproduction in this framework, and the parameters are well constrained by the new data together with the previous πN and γN partial waves [2]. Here we would like to address the indication of a Δ^* resonance around 2000 MeV, namely the $F_{35}(2000)$ state.

In the current version of the compilation of the Particle Data Group [19], the $F_{35}(2000)$ is only rated as a two-star resonance, which means that it has little experimental evidence up to now. Its mass and width are found to be 2160 and 313 MeV in our analysis, respectively. These values should be compared to 2015(24) and 500(52) MeV in the KSU survey [20], the only PWA group presently including this state. Both results find that its major decay channel is the $2\pi N$, whose partial width is bigger than 90%. After adding the $F_{35}(2000)$ into our model, it improves the high energy tail of the partial waves of elastic πN channels. So it seems to be of a non-resonant nature and plays a background role in the πN reactions. This tells us why this state is hard to

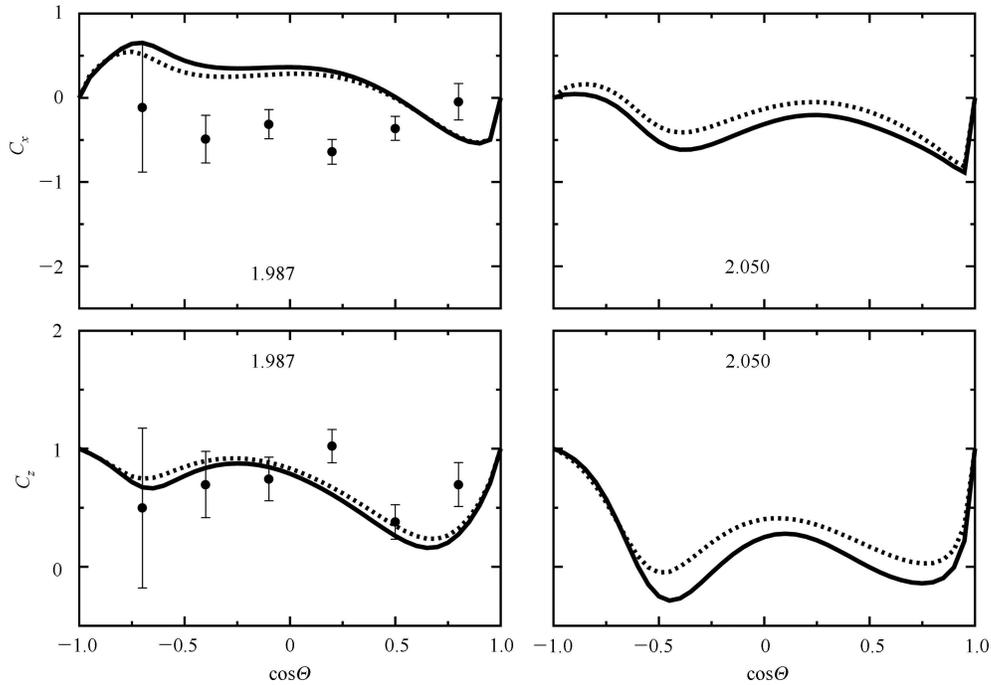


Fig. 1. The spin transfer coefficient C_x and C_z in $\gamma p \rightarrow K^+ \Sigma^0$ reaction. The dashed and solid lines are the full model calculation and the model calculation with the $F_{35}(2000)$ turned off, respectively. Data are taken from CLAS [21]. The numeric values label the center of mass energies in units of GeV.

present clearly in the πN collisions. However, after shutting off this state, the χ^2 increases to a value of 2.3 in the $\gamma p \rightarrow K\Sigma$ channel. Though its contribution to the total cross section is tiny, the role in the polarization observables is clearly seen, as depicted in Fig. 1. Above 2.0 GeV it is shown that the $F_{35}(2000)$ would be more obvious. Other resonances with different quantum numbers, e.g. S_{11} , S_{31} , P_{11} and P_{31} resonance with Breit-Wigner mass varying from 1700 MeV to 2000 MeV are tried to add to the model but we do not find any improvement in the description of the current $K\Sigma$ photoproduction data.

As a matter of fact, the $F_{35}(2000)$ is found to be important in describing the differential cross section of the $\Sigma(1385)$ photoproduction from a proton within a Regge-plus-resonance approach [22]. The determined helicity amplitudes $A_{1/2} = -10 \text{ GeV}^{-1/2}$ and $A_{3/2} = -28 \text{ GeV}^{-1/2}$ should be compared to our fitted values $18 \text{ GeV}^{-1/2}$ and $-23 \text{ GeV}^{-1/2}$, respectively. The corresponding values of the KSU survey [20] are $-16(81) \text{ GeV}^{-1/2}$ and $158(32) \text{ GeV}^{-1/2}$. Though these values have large uncertainties, they consistently indicate the relative big radiative decay width of $F_{35}(2000)$, so the photoproductions are expected to be suitable for studying its properties. In order to resolve unambiguously the signal of $F_{35}(2000)$, it is necessary to extend the coupled-channel model to higher energies, where the $K\Sigma(1385)$, $K^*\Sigma$ and $K^*\Lambda$ channels are opening. It would be also meaningful to explore the $3\pi N$ channel because strong coupling of $F_{35}(2000)$ to $\pi\pi\Delta$ is found [23]. It is meaningful to perform this kind of analysis considering the increasing experimental database in this energy range.

3 New mesonic states in the bottomonium sector

In the energy range below 2.0 GeV, the exotic mesons usually have broad widths and overlap with other non-exotic states. So it is very hard to clearly establish them though a lot of effort has been devoted to the PWA within isobar models [24]. But in the charm and bottom sector, the widths of mesons are limited by flavor changing weak interactions so the possibility of the discovery of exotic states is larger. A typical example is the charged bottomonium-like states, the $Z_b(10610)$ and $Z_b(10650)$, found in the decay of the bottomonium states [25]. Their masses are close to $B^*\bar{B}$ and $B^*\bar{B}^*$ threshold, respectively and the observed decay channel requires a $b\bar{b}d\bar{u}$ quark flavor structure component. If they are confirmed, then we expect other states in the bottom sector in the coupled-channel model, as hinted by the heavy quark flavor symmetry [26]. As a simple illustration, we construct the kernel V from chiral Lagrangians within the framework of the hidden gauge formalism, similar to the case in the charm sector. Then the detailed formalism of ker-

nel V , decomposed in terms of partial waves, could be adopted from Ref. [27–30]. Following this prescription, the propagator of the exchanged meson is proportional to the inverse of its squared mass, so heavy meson exchange is suppressed by a factor $\gamma = (m_L/m_H)^2$, which is also the scale of the $SU(4)$ broken symmetry.

The propagator matrix G can be calculated as a two-meson loop function using dimensional regularization [29]:

$$G_1 = \frac{1}{16\pi^2} \left(\alpha_1 + \text{Log} \frac{m_1^2}{\mu^2} + \frac{M_1^2 - m_1^2 + s}{2s} \text{Log} \frac{M_1^2}{m_1^2} + \frac{p}{\sqrt{s}} \left(\text{Log} \frac{s - M_1^2 + m_1^2 + 2p\sqrt{s}}{-s + M_1^2 - m_1^2 + 2p\sqrt{s}} + \text{Log} \frac{s + M_1^2 - m_1^2 + 2p\sqrt{s}}{-s - M_1^2 + m_1^2 + 2p\sqrt{s}} \right) \right), \quad (3)$$

where M_1 and m_1 are the masses of the intermediate two mesons and p is the corresponding three-momentum in the center of the mass frame. The cut-off μ is set to 2.1 GeV and the α_1 is the subtraction constant. Unitarity is ensured by the equation $\text{Im}(G_1) = -p/8\pi\sqrt{s}$, so the states are present as poles in the complex plane of the full T -matrix in Eq. (1). The two-meson channels are included as:

- 1) Vector-Vector (VV) interaction: $B^*\bar{B}^*$, $B_s^*\bar{B}_s^*$, $K^*\bar{K}^*$, $\rho\rho$, $\omega\omega$, $\phi\phi$, $\omega\phi$, $\Upsilon(1S)\Upsilon(1S)$, $\Upsilon(1S)\omega$, $\Upsilon(1S)\phi$;
- 2) Vector-Pseudoscalar (VP) interaction: $B^*\bar{B}$, $B_s^*\bar{B}_s$, $K^*\bar{K}$, $\rho\pi$, $\phi\eta(')$, $\omega\eta(')$, $\Upsilon(1S)\eta(')$, $\Upsilon(1S)\eta_b$, $\eta_b\omega$, $\eta_b\phi$;
- 3) Pseudoscalar-Pseudoscalar (PP) interaction: $B\bar{B}$, $B_s\bar{B}_s$, $K\bar{K}$, $\pi\pi$, $\eta(')\eta(')$, $\eta_b\eta_b$, $\eta_b\eta(')$.

The masses of the above mesons are adopted from the Particle Data Group [19]. The values of the parameters required to be the input of this coupled-channel model are chosen to be the characterizing values [19, 31]:

1) $m_L = 800.0 \text{ MeV}$ and $m_H = 5370.0 \text{ MeV}$ for the suppression factor γ ;

2) The decay constant: $f_\pi = 93.0 \text{ MeV}$, $f_B = 195.0/\sqrt{2} \text{ MeV}$, $f_{B_s} = 250.0/\sqrt{2} \text{ MeV}$, and $f_{\eta_b} = 801.0/\sqrt{2} \text{ MeV}$.

If the $Z_b(10610)$ and $Z_b(10650)$ are assumed to be dynamically generated states with $I^G(J^{PC}) = 1^+(1^{+-})$, close to $B^*\bar{B}$ and $B^*\bar{B}^*$ thresholds respectively, then the free parameter α_1 could be determined to be -3.12 ± 0.08 for the channels with two heavy vector mesons and -1.65 ± 0.05 for other cases. The uncertainties are estimated from the experimental errors of the masses and widths of the two Z_b states. Then the states of other partial waves could be predicted in Table 1. As can be seen, the states close to the $B^*\bar{B}^*$ threshold are expected to be present in all quantum numbers in the VV channel. These states are very narrow except the $0^-(1^{+-})$ state

with around 200 MeV width. We also predict another narrow $0^-(1^{+-})$ state located at about 10471.1 MeV, but its mass has a relatively large uncertainty. In the PP channel, a $0^+(0^{++})$ state with the pole at around 10732.8 MeV and a $1^-(0^{++})$ state at about 10558.0 MeV are expected, whose widths are both small.

Table 1. Pole positions of non-strangeness states in coupled-channel model with chiral Lagrangians. The inputs are [25]: (a) $Z_b(10650)$ with $M=10652.2\pm 1.5$ MeV and $\Gamma=11.5\pm 2.2$ MeV. (b) $Z_b(10610)$ with $M=10607.2\pm 2.0$ MeV and $\Gamma=18.4\pm 2.4$ MeV.

| | $I^G(J^{PC})$ | RE(\sqrt{s})/MeV | IM(\sqrt{s})/MeV |
|----|---------------|------------------------|----------------------|
| VV | $0^+(0^{++})$ | 10650.2 \pm 1.9 | 11 \pm 7 |
| | $0^-(1^{+-})$ | 10471.1 \pm 42.3 | 1 \pm 1 |
| | | 10650.2 \pm 1.5 | 127 \pm 16 |
| | $0^+(2^{++})$ | 10650.2 \pm 1.7 | 13 \pm 8 |
| | $1^-(0^{++})$ | 10650.2 \pm 1.6 | 7 \pm 5 |
| | $1^+(1^{+-})$ | 10650.2 ^(a) | 5 \pm 4 |
| | $1^-(2^{++})$ | 10650.2 \pm 1.6 | 7 \pm 3 |
| VP | $1^+(1^{+-})$ | 10604.2 ^(b) | -6 \pm 4 |
| | $1^-(1^{++})$ | — | — |
| | $0^+(1^{++})$ | 10604.2 \pm 1.7 | -95 \pm 17 |
| | | 10781.0 \pm 1.3 | -118 \pm 12 |
| | $0^-(1^{+-})$ | 10781.7 \pm 1.3 | -2 \pm 1 |
| PP | $0^+(0^{++})$ | 10732.8 \pm 1.6 | -6 \pm 4 |
| | $1^-(0^{++})$ | 10558.0 \pm 1.1 | -7 \pm 4 |

In the Lippmann-Schwinger equation respecting heavy quark flavor symmetry, a $0^+(1^{++})$ $B^*\bar{B}$ bound state with a mass about 10580 MeV is predicted [26], which should be compared to two broad states with the poles respectively at 10604.2 MeV and 10781.0 MeV in our coupled-channel model. In addition, our model predicts a narrow $0^-(1^{+-})$ state with the pole at 10781.0 MeV. So it is possible to distinguish these two models in the VP channel by experiments.

In Ref. [32] the coupled-channel Schrödinger equation is solved numerically with the meson exchange potentials. They predict several possible bound and resonant states after reproducing the $Z_b(10610)$ and $Z_b(10650)$. A similar framework addresses the role of S-D mixing [33, 34]. These works find that π -meson exchange is important [32–37]. However, in Ref. [26] it is shown that one pion exchange only slightly changes the central value of the calculated mass. In our model, we assume the vector-meson dominance. So it is still an open question as to which interaction, π - or vector-meson exchange, is dominant. The predicted masses of the states in different models would be tested in future experiments.

It is worth mentioning that while the broad states would be completely immersed into the background so as not to be reachable in the high energy machine, the

predicted narrow states could emerge in future experiments. We suggest that the experimentalists search for these narrow states in the photo-induced reactions, i.e. $\gamma N \rightarrow NZ_b$ reactions. The maximum total cross section of charmonium-like mesons in the photo-induced reactions is estimated to be around $\sigma_{Z_c} \lesssim 10$ nb with the assumption of dominance of t -channel vector-meson exchange [38, 39]. If the width of bottomonium-like mesons decaying to $\Upsilon(1S)V$ or $\eta_b V$ channel is at the level of MeV, similar to that of charmonium-like states, the production cross sections of $\gamma N \rightarrow NZ_b$ reactions could be roughly estimated to be at the level:

$$\sigma_{Z_b} = \sigma_{Z_c} \frac{s_c(k_c^\gamma f_{\eta_c})^2}{s_b(k_b^\gamma f_{\eta_b})^2} \lesssim 0.5 \text{ nb}, \quad (4)$$

where s_b and k_b^γ are invariant mass squared and photon energy in the center of the mass system. The factors $s_b(k_b^\gamma)^2$ and $s_c(k_c^\gamma)^2$ could be understood from phase spaces and the decay constants f_{η_b} and $f_{\eta_c} = 420.0/\sqrt{2}$ MeV [29] from the production amplitudes. Though production cross sections are small, the background is clean at this energy range and under control by the Pomeron exchange, leading to the possibility of finding these states, especially the exotic states with 1^{+-} quantum numbers.

4 Remarks and conclusion

The rapid advancement of the coupled-channel model and the wide application of the PWA tools have promoted the recent progress of the hadronic spectra. The coupled-channel model with effective or chiral Lagrangians could be used to predict several hadronic states based on the current empirical information. In this paper we have addressed the evidence of several baryonic and mesonic states with definite quantum numbers in the coupled-channel models. We point out the $K\Sigma(1385)$ photoproduction data would offer a good chance to confirm the $F_{35}(2000)$, a state not yet well established in the current $K\Sigma$ photoproduction data. We have also shown that the coupled-channel model on the basis of the known states could be used to predict the bottomonium-like mesons. Though the production cross section is estimated to be small, the photo-induced reactions in the future high luminosity accelerator seem to have the potential for further study of these states and to disentangle the nature of resonances in different models.

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