# Track segment finding with CGEM-IT and matching to outer drift chamber tracks in the BESIII detector\*

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Abstract: The relative differences in coordinates of Cylindrical Gas Electron Multiplier Detector-based Inner Tracker (CGEM-IT) clusters are studied to search for track segments in CGEM-IT for the BESIII experiment. With the full simulation of single muon track samples, clear patterns are found and parameterized for the correct cluster combinations. The cluster combinations satisfying the patterns are selected as track segment candidates in CGEM-IT with an efficiency higher than 99%. The parameters of the track segments are obtained by a helix fitting. Some  $\chi^2$  quantities, evaluating the differences in track parameters between the track segments in CGEM-IT and the tracks found in the outer drift chamber, are calculated and used to match them. Proper  $\chi^2$  requirements are determined as a function of transverse momentum and the matching efficiency is found to be reasonable.

**Keywords:** CGEM-IT, track segment finding, track matching

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## 1 Introduction

The inner drift chamber (IDC) of the Beijing Spectrometer III (BESIII) [1] detector has shown a gradual aging effect. The hit efficiency and spatial resolution will deteriorate and will affect the reconstruction of charged tracks in the future. A Cylindrical Gas Electron Multiplier Detector-based [2] Inner Tracker (CGEM-IT) has been proposed as a candidate for the upgrade of the IDC [3].

The proposed CGEM-IT consists of three layers of CGEM detector. Each CGEM comprises one cathode, three GEM foils and one anode. There are two kinds of strips on the anode readout plane [4] which point in two different directions and are used to measure the induced charge from charged particles. The GEANT4-based [5] full simulation package was implemented for CGEM-IT, including a detailed detector description and a simplified digitization [6]. A cluster reconstruction algorithm was developed, in which the continuous fired strips are reconstructed as clusters and the positions are calculated with the charge centroid method [6]. A track fitting algorithm for the drift chamber, based on the Kalman Fil-

ter method [7], was also extended to incorporate both CGEM-IT clusters and outer drift chamber (ODC) hits, so that the spatial and momentum resolution can be studied for the tracks reconstructed within the tracking system of the CGEM-IT and the ODC [6]. This paper describes the track segment finding with CGEM-IT and the track matching between CGEM-IT and ODC.

# 2 Track segment finding with CGEM-IT

As the CGEM-IT includes three layers, usually three clusters can be reconstructed when a charged particle goes through the CGEM-IT in a magnetic field. The two coordinates of the reconstructed clusters are the azimuthal angle  $\phi$  and the position Z along the beam direction, where the center of the detector is taken as the origin point. Figure 1 shows an example of a charged track and the induced three clusters on the transverse plane. The selection of the correct three-cluster combinations is called track segment finding with CGEM-IT. The relation of the three clusters from one track can be studied with the relative differences in  $\phi$  and Z between the first or the third layer and the second layer.

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Then each combination of three clusters can be shown as one point on a two-dimensional plane of  $\delta\phi$ , where the two dimensions are  $\delta\phi_{21}=\phi_1-\phi_2$  and  $\delta\phi_{23}=\phi_3-\phi_2$  (the subscript indexes 1, 2, 3 indicate the layer number). Similarly the combinations of three clusters can also be shown on a two-dimensional plane in  $\delta Z$  ( $\delta Z_{21}=Z_1-Z_2$ ,  $\delta Z_{23}=Z_3-Z_2$ ).

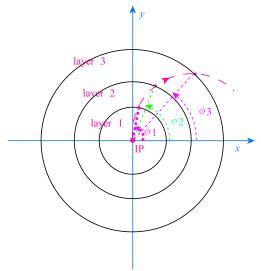


Fig. 1. (color online) The projection on the transverse plane (perpendicular to beam direction) of the three layers (black circles) of CGEM, and a charged track (red dashed line). The three intersections between the CGEM layers and the track are called clusters, the azimuth angles of which are  $\phi_1$ ,  $\phi_2$  and  $\phi_3$  respectively.

#### 2.1 Pattern in $\delta \phi$

Based on the full simulation, the correct threecluster-combinations on the  $\delta\phi$  plane for  $\mu^{\pm}$  with a uniformly distributed transverse momentum  $p_t$  between 0.05 and 1.0 GeV/ $c^{-1}$ , and with a polar angle  $\theta$  satisfying  $|\cos\theta| < 0.93$ , are shown in Fig. 2. A clear anti-diagonal pattern can be observed, the general behavior of which is fitted to a linear function through the origin point. As shown in Fig. 3, the scattering points on the  $\delta\phi$  plane can be presented by a signed perpendicular distance  $d_{\delta\phi}$ to the fitted line (negative for points bellow the line, positive for points above the line) and the projected position  $\delta L_{\delta\phi}$  on the line. The  $d_{\delta\phi}$  distributions at different  $\delta L_{\delta\phi}$ are studied and fitted to a Gaussian function to extract the mean value  $\mu_{\delta\phi}$  and the resolution  $\sigma_{\delta\phi}$ . One example of the  $d_{\delta\phi}$  distribution for  $0.054 < \delta L_{\delta\phi} < 0.108$  rad is shown in Fig. 4. The obtained mean and resolution values at different  $\delta L_{\delta\phi}$  are shown in Fig. 5 and can be fitted to a third order polynomial function without even terms and a fourth order polynomial function without

odd terms respectively. According to the parameterized pattern, the central line of the pattern and a selection window corresponding to a range of  $\pm 3\sigma_{\delta\phi}$  for the track segment candidates on the two dimensional  $\delta\phi$  plane are determined and shown in Fig. 2.

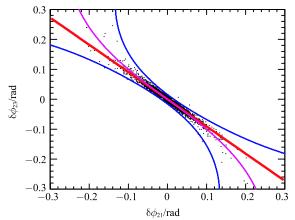


Fig. 2. (color online) The distribution of correct three-cluster-combinations (black points) on the  $\delta\phi$  plane for  $\mu^{\pm}$  with random  $p_{\rm t}$  ranging from 0.05 to 1.0 GeV/c. The red line is the fitted result of the distribution to a linear function through the origin point. The purple curve is the central line of the pattern found, which is obtained by fitting the residual distributions along the straight line. The area between the two blue curves is the track segment selection window on the  $\delta\phi$  plane which corresponds to a range of  $\pm 3\sigma_{\delta\phi}$ .

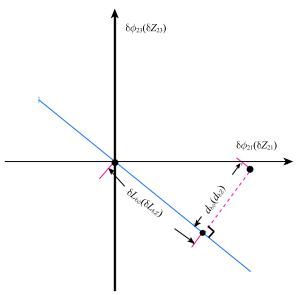


Fig. 3. The definitions of the signed perpendicular distance  $d_{\delta\phi}(d_{\delta Z})$  and the projected position  $\delta L_{\delta\phi}(\delta L_{\delta Z})$ . The black dot is a certain three-cluster combination and the line is the fitted result described in the text.

<sup>&</sup>lt;sup>1)</sup>The behavior of higher momentum charged tracks is similar, so the momentum of the charged tracks under study are limited up to 1 GeV/c.

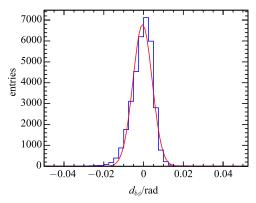


Fig. 4. The  $d_{\delta\phi}$  distribution for  $0.054 < \delta L_{\delta\phi} < 0.108$  rad. The curve is the fitted result to a Gaussian function.

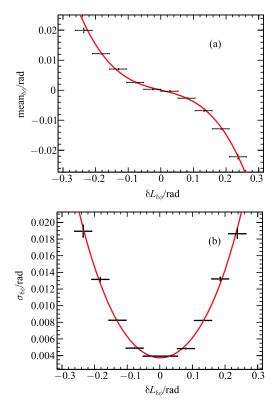


Fig. 5. (color online) The mean (a) of  $d_{\delta\phi}$  as a function of  $\delta L_{\delta\phi}$  on the  $\delta\phi$  plane is fitted to a third order polynomial without even terms (red line). The resolution  $\sigma_{\delta\phi}$  (b) of  $d_{\delta\phi}$  as a function of  $\delta L_{\delta\phi}$  is fitted to a fourth order polynomial without odd terms (red line).

#### 2.2 Pattern in $\delta Z$

Correct three-cluster combinations on the  $\delta Z$  plane are studied at different transverse momenta with the same simulated  $\mu^{\pm}$  sample. Figure 6 shows an example for  $\mu^{\pm}$  with  $p_{\rm t}$  between 0.15 and 0.25 GeV/c and a nice anti-diagonal pattern emerges which can also be fitted to a straight line going through the origin. As shown

in Fig. 7, the slope of the fitted line has a slight transverse momentum  $p_t$  dependence. Similarly, the signed distance  $d_{\delta Z}$  to the linear line and projected position  $\delta L_{\delta Z}$  for three-cluster combinations on the  $\delta Z$  plane are also defined as shown in Fig. 3, and the mean and resolution of  $d_{\delta Z}$  distributions at different  $\delta L_{\delta Z}$  and  $p_t$  are studied by fits to a Gaussian function (one example is shown in Fig. 8). The obtained mean values of  $d_{\delta Z}$  on the  $\delta Z$  plane are consistent with zero. The resolution of  $d_{\delta Z}$  on the  $\delta Z$  plane for a certain  $p_t$  can be fitted to a second order polynomial function with even terms only as shown in Fig. 9.

For  $\mu^{\pm}$  with a certain  $p_t$  interval, the  $\pm 3\sigma_{\delta Z}$  area for three-cluster-combination selection can be determined with the corresponding slope and resolution. But before the track reconstruction, the momentum of charged particles is unknown. So the  $3\sigma_{\delta Z}$  areas for different  $p_t$  are merged together to define a selection window for charged tracks with  $p_t > 0.05$  GeV/c, which is shown in Fig. 6.

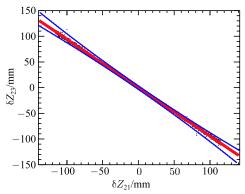


Fig. 6. (color online) The distribution of correct three-cluster combinations (black points) on the  $\delta Z$  plane for  $\mu^\pm$  with  $p_{\rm t}$  between 0.15 and 0.25 GeV/c. The red line is the fitted result of the distribution to a linear function through the origin point. The area between the two blue lines is the merged three-cluster-combination selection window on the  $\delta Z$  plane as described in the text.

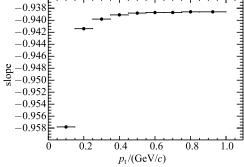


Fig. 7. The slope of the fitted straight line on the  $\delta Z$  plane as a function of  $p_{\rm t}$  for the correct three-cluster combinations from full simulated  $\mu^{\pm}$  samples.

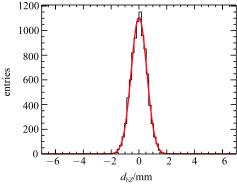


Fig. 8. The  $d_{\delta Z}$  distribution for  $0 < \delta L_{\delta Z} < 34.3$  mm. The curve is the fitted result to a Gaussian function.

# 2.3 Efficiency of track segment finding

With the defined selection windows, as described in the previous subsections, on both of the  $\delta\phi$  and  $\delta Z$  planes, the three-cluster combinations are selected as the track segment candidates in CGEM-IT. All the three-cluster combinations and those after the selection are shown in Fig. 10 for an independent full simulated  $\mu^{\pm}$  samples. The three-cluster combinations inconsistent with the found patterns are due to the clusters induced by electrons from muon decay. The efficiency of the track

segment finding is computed as the ratio of the number of correct three-cluster combinations after selection over the number of generated tracks. The obtained efficiency of the track segment finding as a function of  $p_t$  for simulated  $\mu^{\pm}$  is shown in Fig. 11, and ranges from 99.1% to 99.7% depending on  $p_t$ . The lower efficiency at lower  $p_t$  is caused by the decay of muons.

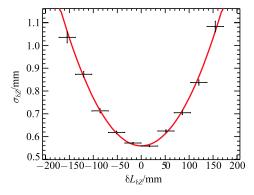


Fig. 9. The resolution  $\sigma_{\delta Z}$  of  $d_{\delta Z}$  at different  $\delta L_{\delta Z}$  for  $\mu^{\pm}$  with  $p_{\rm t}$  between 0.15 and 0.25 GeV/c (black dots with error bar). The curve is the fitted result to a second order polynomial with even terms only.

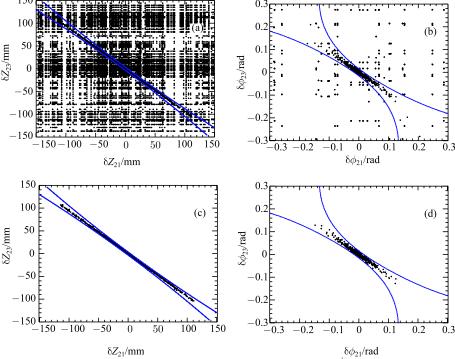


Fig. 10. (color online) The distributions of three-cluster combinations for  $\mu^{\pm}$  with  $p_{\rm t} > 0.05~{\rm GeV/c}$ . Before the selection, three-cluster combinations are shown on the  $\delta Z$  (a) and  $\delta \phi$  (b) plane. Most correct three-cluster combinations are selected by both selection windows on the  $\delta Z$  and  $\delta \phi$  plane, and their distribution are (c) and (d). The backgrounds are incorrect three-cluster combinations in which some clusters are produced by electrons from muon decay.

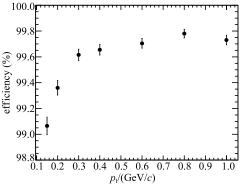


Fig. 11. The efficiency of the track segment finding as a function of  $p_t$  for full simulated  $\mu^{\pm}$  samples.

#### 2.4 Track segment fitting

To describe quantitatively the found track segments in CGEM-IT, the three-cluster combinations are fitted to a helix model with five parameters. Some of the parameters are defined with respect to a referenced three-dimensional point called the pivot, which can be arbitrary but usually is the interaction point or the origin point. The five parameters are

dr: the signed distance on the transverse plane from the pivot to the helix,

 $\phi_0$ : the azimuthal angle specifying the pivot with respect to the helix center,

 $\kappa$ : the product of the particle charge and the reciprocal of  $p_t$ ,

 $\mathrm{d}z$ : the signed distance of the helix from the pivot in the Z direction,

 $\tan \lambda$ : the slope of the track and  $\lambda = \pi/2 - \theta$  where  $\theta$  is the polar angle of the track.

The least squares method is used for the helix fitting, and the  $\chi^2$  for the fitting is defined as

$$\chi_{\text{fit}}^2 = \Sigma_j \Sigma_i \frac{(X_{i,j}(H) - X_{i,j,\text{measured}})^2}{\sigma_{i,j}^2}, \tag{1}$$

where X(H) is the cluster position calculated with the helix parameters H, and  $X_{\text{measured}}$  means the cluster position measured by the CGEM-IT.  $\sigma$ , which comes from

the residual between the measured result and MC truth information is the error of the measured cluster positions,  $i=1,\ 2$  stand for  $\phi$  and Z respectively, and j=1,2,3 stand for the 3 clusters. The  $\chi^2_{\rm fit}$  minimization is realized with MINUIT [10].

Since the high  $p_{\rm t}$  track is quite straight in CGEM-IT, only three clusters close by with limited spatial resolution are unable to determine  $\kappa$  well. The effect is reflected in the success rate of the 5-parameter helix fitting which drops at high  $p_{\rm t}$  as shown in Fig. 12 (black dots). To improve the success rate of the helix fitting for high  $p_{\rm t}$  tracks,  $\kappa$  can be fixed to the one calculated with the three transverse positions of the clusters if the 5-parameter helix fitting fails. The success rate of the helix fitting after fixing  $\kappa$  is kept above 99.4% and reaches 99.7% at high  $p_{\rm t}$  as shown in Fig. 12 (blue triangles).

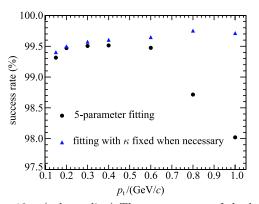
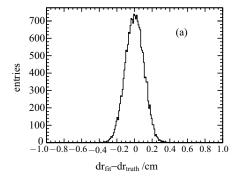
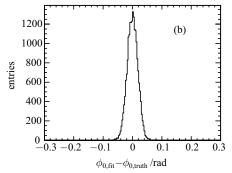


Fig. 12. (color online) The success rate of the helix fitting as a function of  $p_t$  with all 5 parameters floating (black dots) and with  $\kappa$  fixed if the previous 5-parameter fitting fails (blue triangles).

The residual distributions for the five helix parameters after the track segment fitting are shown in Fig. 13, where the residual is defined by  $\delta H_i = H_i^{\rm fit} - H_i^{\rm truth}$  (i = 1-5), namely the difference between the fitted helix parameters  $H_i^{\rm fit}$  and the true helix parameters  $H_i^{\rm truth}$  at the  $\mu^{\pm}$  track generation. These distributions demonstrate that reasonably unbiased track parameters can be obtained by the fitting.





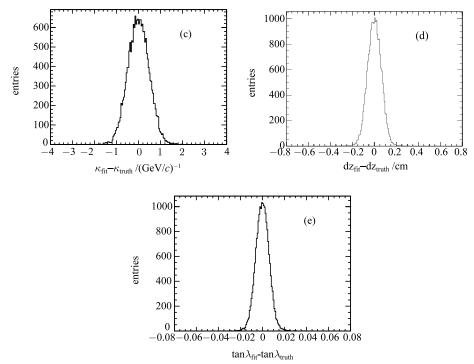


Fig. 13. The distributions of residuals for the helix parameter dr (a),  $\phi_0$  (b),  $\kappa$  (c), dz (d) and  $\tan \lambda$  (e), which are obtained from a full simulated  $\mu^{\pm}$  with  $p_t = 0.8 \text{ GeV}/c$ .

# 3 Matching CGEM-IT track segments to ODC tracks

Typically, a complete track is found if both the track segment in CGEM-IT and the track in the ODC are found and matched. The tracks in the ODC can be reconstructed by the track finding algorithm [8, 9], the helix parameters of which are also obtained by fitting. To evaluate the consistency between the track segments in CGEM-IT and the tracks in the ODC, the quantity  $\chi^2_{\text{match}}$  is calculated for each of the 5 helix parameters:

$$\chi_{\text{match}}^2 = \frac{(H_{\text{CGEM-IT}} - H_{\text{ODC}})^2}{\sigma_{\text{CGEM-IT}}^2 + \sigma_{\text{ODC}}^2},$$
 (2)

where H is a parameter of the helix from the track segment in CGEM-IT or the track in the ODC, and  $\sigma$  is the corresponding error. A global  $\chi^2$  of the five track parameters can be defined. But the error matrices are not perfectly evaluated here, as some effects can not be considered in a simple helix fitting (for instance, the inhomogeneity of the magnetic field, energy loss of charged tracks and so on). This affects the five parameters unequally, so it is better to study the parameters independently. As an example, the  $\chi^2_{\rm match}$  distributions of the 5 parameters for a simulated  $\mu^{\pm}$  sample with  $p_t = 0.8 \ {\rm GeV}/c$  are shown in Fig. 14. Appropriate  $\chi^2_{\rm cut}$  values are chosen to ensure that the individual efficiency for each parameter is not less than 99.5% when  $\chi^2_{\rm match} < \chi^2_{\rm cut}$  is required. The determined  $\chi^2_{\rm cut}$  values

at different  $p_t$  for the 5 helix parameters are shown in Fig. 15.

The  $\chi^2_{\rm cut}$  value at a  $p_{\rm t}$  between two neighboring points on the plot is obtained by a linear interpolation. If each of the five  $\chi^2_{\rm match}$  is less than the corresponding  $\chi^2_{\rm cut}$  value, the track segment in the CGEM-IT and the track in ODC are called matched.

In the study described here, the procedure of finding a complete track, in the tracking system consisting of the CGEM-IT and the ODC, includes the track segment finding in CGEM-IT (black dots), track segment fitting (green inverted triangles), ODC track finding (red upright triangles) and the track matching (blue stars). These steps proceed in a sequence. The absolute efficiency after each step is shown in Fig. 16 at different  $p_t$ . The relative efficiency for the track matching is around 99% and the total efficiency after the track matching ranges from 94.6% to 97.0% with some  $p_t$  dependence.

# 4 Summary and outlook

As one candidate for the upgrade of the IDC at BE-SIII, the CGEM-IT has been proposed and designed. The implementation and studies of the full simulation of the CGEM-IT and the track reconstruction are in progress. With a full simulated  $\mu^{\pm}$  sample, two patterns are found in the relative difference in azimuth angle and in the position along the beam direction. These pat-

terns are parameterized and used to select the cluster-combinations as track segments in CGEM-IT with an efficiency more than 99%. The quantitative description of the track segment is obtained by a helix fitting. The chi-squared quantities evaluating the consistency between the track segments in CGEM-IT and the tracks in the ODC are calculated and used in the track matching. Ap-

propriate requirements on the chi-squares are chosen and the relative efficiency for the track matching is about 99%. A reasonable total efficiency, which is between 94.6% and 97% depending on the transverse momentum of tracks, is obtained after the track segment finding with CGEM-IT, track segment fitting, track finding with ODC and the track matching.

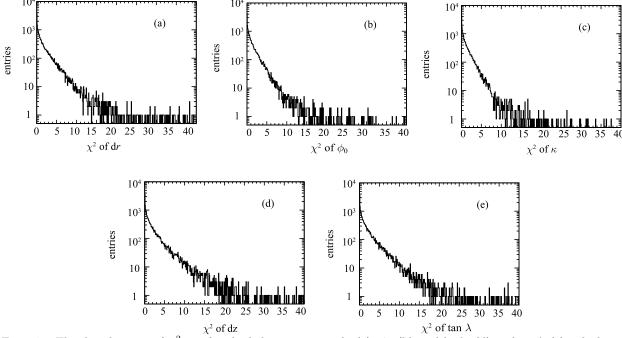


Fig. 14. The distributions of  $\chi^2_{\rm match}$  for the helix parameter dr (a),  $\phi_0$  (b),  $\kappa$  (c), dz (d) and  $\tan \lambda$  (e), which are obtained from a full simulated  $\mu^{\pm}$  sample with  $p_t = 0.8 \ {\rm GeV}/c$ .

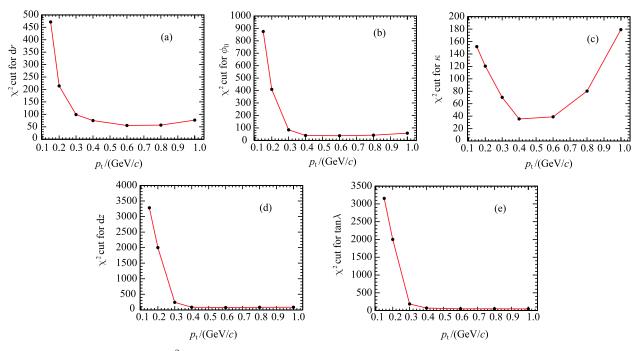


Fig. 15. The  $\chi_{\text{cut}}^2$  values at different  $p_t$  for dr (a),  $\phi_0$  (b),  $\kappa$  (c), dz (d) and  $\tan \lambda$  (e).

To further improve the total track finding efficiency with the tracking system consisting of the CGEM-IT and the ODC, one should optimize the track finding with the ODC and develop a global track finding algorithm with both CGEM clusters and ODC hits when it is hard to find good track segments in CGEM-IT or tracks in ODC. Additional efforts should be devoted to the tracking of low momentum charged particles. All the charged track reconstruction algorithms should also be fine tuned for different particles (e $^{\pm}$ ,  $\mu^{\pm}$ ,  $\pi^{\pm}$ ,  $K^{\pm}$ ,  $p/\bar{p}$ ) with the existence of backgrounds.

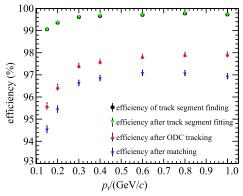


Fig. 16. The absolute efficiency at different  $p_t$  after the track segment finding in CGEM-IT, track segment fitting, ODC track finding and the track matching.

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